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## Warehouse Commodity Classification from Fundamental Principles. Part II: Flame Heights and Flame Spread

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### Abstract

In warehouse storage applications, it is important to classify the burning behavior of commodities and rank them according to their material flammability for early fire detection and suppression operations. In this study, an approach is presented to model the large-scale effects of warehouse fires by decoupling the problem into separate processes of heat and mass transfer. Two nondimensional parameters are proposed to represent the physical phenomena at the large-scale, a mass transfer number and the soot yield of the fuel that controls the radiation observed in the large-scale. In order to facilitate modeling, a mass transfer number (or B-number) was experimentally obtained using mass-loss (burning rate) measurements from bench-scale tests, following from a procedure that was developed in Part I of this paper.

Two fuels are considered: corrugated cardboard and polystyrene. Corrugated cardboard provides a source of flaming combustion in a warehouse and is usually the first item to ignite and sustain flame spread. Polystyrene is typically used as the most hazardous product in large-scale fire testing. A mixed fuel sample (corrugated cardboard backed by polystyrene) was also tested to assess the feasibility of ranking mixed commodities using the bench-scale test method. The nondimensional mass transfer number was then used to model the upward flame propagation on 6.1 – 9.1 m (20 – 30 ft) stacks of corrugated

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cardboard boxes on rack-storage during the initial period of flame spread (involving flame spread over the corrugated cardboard face only). Good agreement was observed between the model and large-scale experiments during the initial stages of fire growth.

*Keywords:*

upward flame spread, flame height, commodity classification, B number, Group A plastic, warehouse fire

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## Nomenclature

### Symbols

$A_l$	Eq. 10a
$A_t$	Eq. 10b
$B$	B-number / Spalding Mass Transfer Number, Eq. (2) (–)
$c_p$	Specific heat ( $J/g\cdot K$ )
$D$	Species diffusivity ( $m^2/s$ )
$d$	Panel separation distance (m)
$Gr$	Grashof number (–)
$\Delta H_c$	Heat of combustion ( $J/g$ )
$\Delta H_g$	Heat of gasification ( $J/g$ )
$g$	Acceleration due to gravity ( $m/s^2$ )
$h$	Heat transfer coefficient ( $W/m^2\cdot K$ )
$h_c$	Convective heat transfer coefficient ( $W/m^2\cdot K$ )
$h_r$	Radiant heat transfer coefficient ( $W/m^2\cdot K$ )
$k$	Thermal conductivity ( $W/m\cdot K$ )
$\dot{m}''$	Mass-loss rate per unit area ( $g/m^2\cdot s$ )
$Nu$	Nusselt number (–)
$Pr$	Prandtl number (–)
$Q$	Energy losses at fuel surface (W)
$\dot{q}_A''$	Volumetric heat release rate ( $kW/m^3$ )
$\dot{q}_c''$	Convective heat flux per unit area ( $kW/m^2$ )
$\dot{q}_{loss}''$	Surface heat loss rate ( $kW/m^2$ )
$\dot{q}_r''$	Radiant heat flux per unit area ( $kW/m^2$ )
$\dot{q}''(x)$	Surface heat flux per unit area ( $kW/m^2$ )
$\dot{q}''(0)$	Surface heat flux at pyrolysis height ( $kW/m^2$ )
$\dot{q}_F'$	Rate of forward heat transfer per unit width (W/m)
$\dot{q}_c'$	Rate of heat release by combustion per unit width (W/m)
$r$	Mass consumption number ( $Y_{O,\infty}/\nu_s$ )
$T_m$	Average temperature between flame and fuel surface (K)
$T_p$	Fuel pyrolysis temperature (K)
$T_\infty$	Ambient temperature (K)
$v_s$	Stoichiometric oxygen-mass fuel ratio (–)
$U$	Free stream velocity (K)
$V_{xp}$	Velocity of pyrolysis front (m/s)
$w$	Panel/sample width (m)
$x_f$	Flame height (m)

$x_p$	Pyrolysis height ( $m$ )
$Y_{O_2}$	Mass fraction of oxygen ( $g/g$ )
$Y_g$	Soot yield of combustion gases ( $g/g$ )
$Y_s$	Soot yield of fuel ( $g/g$ )

### **Greek Symbols**

$\alpha$	Thermal diffusivity ( $m^2/s$ )
$\beta$	Thermal expansion coefficient ( $1/K$ )
$\chi$	Fraction of flame radiation lost to the environment ( $-$ )
$\delta$	Preheat distance ( $m$ )
$\epsilon$	Emissivity ( $-$ )
$\rho$	Density ( $g/m^3$ )
$\mu$	Viscosity ( $kg/m-s$ )
$\nu$	Kinematic viscosity ( $m^2/s$ )
$\Phi$	Forward heating parameter ( $\Phi = \dot{q}'_F/\dot{q}'_c$ )
$\sigma$	Stefan-Boltzmann Constant ( $W/m^2-K^4$ )
$\tau$	Shear stress at surface ( $Pa$ )
$\varsigma_f$	Nondimensional flame height ( $-$ )
$\varsigma_p$	Nondimensional panel height ( $-$ )

### **Subscripts**

F	Fuel
f	Flame
g	Gas
m	Mean
s	Solid
$\infty$	Ambient

## 1. Introduction

Automatic sprinkler system designs for warehouse storage occupancies are currently based upon a commodity classification system, which is widely acknowledged to involve a subjective design process. Yet, this is marginalized by reference to what is characterized as an excellent record of fire loss history. The question of whether the loss history justifies marginalizing the role of engineers is debatable, but one thing is clear: this history is not relevant in the uncharted territory of large warehouses with storage heights up to 30 m (100 ft). Over the last 50 years, fire protection engineers have relied on large-scale tests to classify commodities into one of seven classes [1] that are representative of their fire performance under specific geometric configurations and ignition conditions. This classification process, which relies on expensive full-scale testing, results in increased safety gaps as the industry creates new and untested materials that are stored in large quantities. Additionally, a categorical classification system places an artificial ceiling on the level of hazard in which a Group A commodity is assumed to be the worst-case fuel load that is possible. Only a limited amount of fundamental science has been performed in this area, which is largely due to the range of complexities that occur in large-scale fire phenomena. Currently, no tests that are known to the authors provide a complete set of fundamental, nondimensional parameters that can be used in engineering calculations towards the safer design of large storage facilities. Efforts that result in the development of such test methods and classification methodology with a sound scientific basis may fulfill an urgent need to improve upon the current warehouse design methods.

The motivation for this study was a series of recent losses in large warehouse storage facilities. Of these incidents, which were reviewed in Part I of this paper, some of the facilities included automatic sprinkler systems that were installed in full compliance with their respective current codes and standards [2]. The negative impacts of these devastating fire incidents were felt by the occupants, firefighters, insurance interests, and local environments. From a business aspect, millions of dollars of materials or products are lost, and operations may be halted [3]. Furthermore, insurance premiums are increased as a result of the fire, and the lost time can never be recovered. From a life-safety aspect, the lives of workers and responding firefighters are endangered, which can result in injuries or death. The water runoff from firefighting operations and the resulting smoke plumes can also adversely affect the sur-

rounding environment. The development of an approach to protect these facilities based upon the combustible materials that are stored, the layout of these materials, and the complex interaction with potential suppression systems is a critical step towards reducing the severity of devastating warehouse losses. As a first step towards improving current protection measures, an improvement of the methods for commodity classification, which ranks the combustibility of one group of stored materials against another, has been assessed.

In Part I of this paper, a method was developed to experimentally quantify the burning rate of a material based upon the nondimensional comparison of a materials chemical energy released during the combustion process with the energy required to vaporize the fuel, which was measured as a B-number. Commodities are classified to design sprinkler protection systems for most warehouse scenarios, and because such a sprinkler systems goal is to suppress or control a fire, the ranking of materials based upon the burning and spread rates of a potential fire is appropriate. Experiments were performed on a standard warehouse commodity, a Group A plastic, which is typically used to represent the worst-case commodity in large-scale tests. The commodity consisted of a single corrugated cardboard box that measured 53 x 53 x 51 cm and contained 125 crystallized polystyrene cups that were segregated by corrugated cardboard dividers. All of the faces except for the front face of the commodity were uniformly insulated, and the front face of the commodity was ignited at its base.

The experimental observations of the Group A plastic commodity resulted in a qualitative description of the burning process over three distinct stages of burning. The first stage was characterized by upward flame spread over the front face of the corrugated cardboard, followed by a decreased burning rate as the cardboard smoldered and the polystyrene heated, and finally a sharp increase in the burning rate after ignition of the polystyrene. Despite the complex configuration, each stage resulted in distinct material involvement, which indicates the potential to model distinct material involvement from each stage using bench-scale testing. Fluctuations between the repeated tests also indicated the difficulty in obtaining repeatable measurements during these larger tests; therefore, small-scale test methods that can be repeated at a level of statistical accuracy may greatly improve the applicability of the results.

Part II of this study continues the development of a nondimensional approach to characterizing the burning behavior of materials. The bench-

scale tests that were performed in this study involved a small, flat sample (5 x 20 cm) of corrugated cardboard or polystyrene oriented vertically in which the burning was isolated to the front surface of the sample. The flow was considered to be laminar due to the observed behavior of the flow in the experiments. At the reduced scale, the tests captured the effects of the commodity material properties on the flame spread process while separating the large-scale effects such as turbulence and radiation. Nondimensional B-numbers were experimentally determined for the samples with greater accuracy than previous experiments. A flame spread model was then utilized to demonstrate the application of the experimentally measured B-numbers to predict flame heights in large-scale configurations. A particular configuration considered in this study is upward flame spread in the flue space between corrugated cardboard, which is typical of warehouse storage arrangements. The model was extended to account for both convective and radiative heat transfer by incorporating convective and radiative heat transfer correlations. This segregated approach captures the condensed phase pyrolysis phenomena by using a nondimensional parameter to represent the mass transfer processes, the gas phase heat transfer by including an appropriate convective heat transfer correlation, and radiative heat transfer effects that are based on previous studies.

## 2. Literature Review

Previous studies have attempted to model some of the large-scale effects of warehouse fires by measuring the relevant parameters using small-scale test methods. One such effort by Hamins and McGrattan [4] constructed single-cell replicates of a Group A plastic commodity. The purpose of the Group A plastic tests was to provide input parameters into a computational fluid dynamics model using a measured heat release rate as the thermal loading input for a large-scale warehouse fire. The model predictions were unable to describe the detailed fire growth in storage applications.

Several studies have addressed the issue of upward flame spread on corrugated cardboard surfaces. Grant and Drysdale [5] modeled the flame spread along corrugated cardboard during the early growth stages of a warehouse fire by adapting the linearized Satio, Quintiere, and Williams [6] flame spread model with Karlsson's [7] burnout length and solving numerically. Dimensional parameters that were obtained experimentally were used as inputs to numerically model the flame height, velocity of the flame front, and pyrolysis

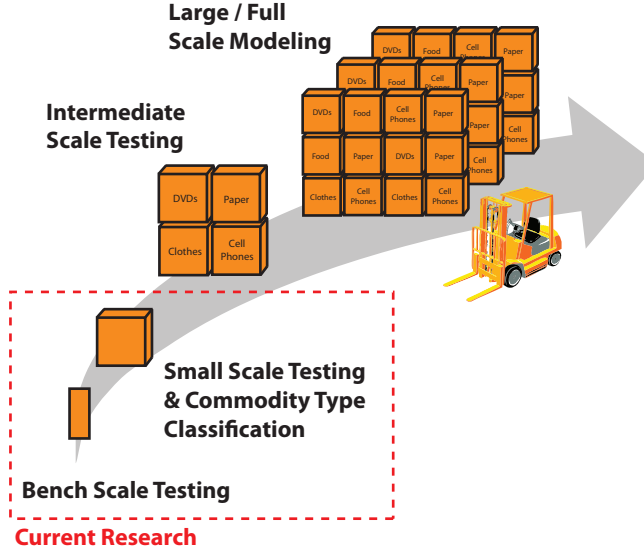


Figure 1: A research approach to the warehouse fire problem. The scales studied in this work are shown by the dashed box.

front progression as a two-dimensional problem. Good agreement between the experimental results and the numerical results were obtained, although the model was found to be sensitive to averaged input parameters, such as the forward heat flux from the flame. Alvares et al. [8] studied the effects of panel separation on vertical flame spread and mass-loss rates in small-scale corrugated cardboard tests to determine the rate of fire growth along vertical flues in warehouses.

Continued efforts by Inganson and de Ris [9] and Inganson [10] have identified the importance of the commodity configuration, the mode of heat transfer, and the flue spacing of commodity boxes in warehouse fires. Inganson's work [9] identified some of the dominant factors in the large-scale warehouse fire growth process, and emphasized the importance of separating the material properties of the fuel from the heat transfer and flow conditions that can result due to the various configurations of the fuel packages. In separating the warehouse fire problem into two distinct phenomena, it then becomes a problem of defining the material properties (condensed phase), flow conditions (geometry), and heat transfer (gas phase). Work performed by de Ris and Orloff [11], de Ris et al. [12], Foley [13], and Foley and Drysdale [14] served to characterize the mode of heat transfer from an upward



propagating flame in a warehouse configuration and to quantify the convective and radiative heat transfer that drives the upward flame spread process in the gas phase. Variations in the heat transfer from the small-scale to the large-scale was shown by de Ris et al. [12] to be related by similarity effects that are present in buoyant, turbulent boundary layer flows. This result can be used to extend the analytical results that were developed for heat and mass transfer in laminar boundary layers to turbulent boundary layers.

The primary driving force of upward flame spread is the heat flux from the advancing fire plume towards the unburned material [15]. In a warehouse setting, this heat transfer may manifest itself in the form of strong convective currents tunneled through flue spaces or as soot-induced radiation from large, luminous fire plumes. As a first approximation, the convective heat transfer can be approximated from the standard engineering correlations for turbulent boundary layer flows. An approach to modeling the radiative heat transfer from the fire plume can be used by adapting an approach from Markstein and de Ris [12], which quantifies the soot yield of a fuel based on its smoke point.

In the early stages of a warehouse fire, before the fire sprinklers are activated, the mass transfer is intrinsically coupled to the material properties of the stored commodity, packing material, and outer corrugated cardboard covering. Due to the different burning behavior of each material, which is also a function of the packing and orientation, the problem of classifying a commodity based on its fire hazard is a complex one. A general approach for describing the heat, mass, and momentum transfer by way of differential equations for simple geometries such as a droplet, flat horizontal, and vertical plate are discussed extensively in previous fire literature [16–19]. Physically, all of these theories rely on the extended Reynolds analogy that includes the combustion of solid fuels [20] in the form

$$\frac{\tau}{U\nu^{2/3}} = \frac{h}{c_p\alpha^{2/3}} = \frac{\dot{m}''}{D^{2/3} \cdot \ln(1 + B)}. \quad (1)$$

Equation 1 is also referred to as the Chilton-Colburn [21] extension to the Reynolds analogy because it incorporates both of the turbulent and laminar molecular processes of diffusion by using the kinematic viscosity or momentum diffusivity ( $\nu$ ), the thermal diffusivity ( $\alpha$ ), and the species diffusivity ( $D$ ). Equation 1 implies that the shear stress at the surface ( $\tau$ ) is related to the heat transfer ( $h/c_p$ ) and mass transfer from combustion ( $\dot{m}''$ ). The terms  $U$ ,  $h$ , and  $c_p$  are the free stream velocity, heat transfer coefficient, and

specific heat of the gas, respectively. The term  $B$  that appears in Eq. 1 is a nondimensional proportionality constant that relates the rate of mass transfer (e.g., vaporization, combustion) to the heat transfer and shear stress, which is referred to as the B-number in this work. A recent study by Raghavan et al. [22] further analyzed this proportionality and showed that Eq. 1 is valid except during the ignition and extinction conditions. Because the B-number in Eq. 1 represents the driving force for mass transfer, it is also referred as the “transfer number” by Spalding [23] and can be represented as the ratio

$$B = \frac{(1 - \chi)Y_{O_2,\infty}(\Delta H_c/r) - c_p(T_p - T_\infty)}{\Delta H_g + Q}, \quad (2)$$

where  $\chi$  is the fraction of radiation lost to the environment,  $Y_{O_2,\infty}$  is the mass fraction of oxygen in the air,  $\Delta H_c$  is the heat of combustion,  $r$  is the mass consumption number given by  $(Y_{O,\infty}/\nu_s)$ ,  $c_p$  is the specific heat of air,  $T_p$  is the vaporization temperature of the fuel,  $T_\infty$  is the ambient temperature,  $\Delta H_g$  is the heat of gasification, and  $Q$  represents the energy losses at the fuel surface [24].

The B-number is composed of material-related properties; therefore, it has been used to rank material flammability in previous fire literature [25–28]. A problem that often remained in the past studies was implementing the B-number to assess large-scale behavior, which will be further discussed in this study. Figure 2 shows the B-numbers for a range of fuels [27] as a function of the pyrolysis temperature of the materials. The circles show the values of the thermodynamic B-number versus the pyrolysis temperatures for the fuels as calculated by Annamalai and Sibulkin [27]. The thermodynamic values of the B-numbers were calculated using Eq. 2 where  $\chi$  and  $Q$  are assumed to be equal to zero, which represents an ideal value with no losses. The liquid fuels shown in Fig. 2 have a larger B-number value and a lower pyrolysis temperature, which corresponds to a smaller amount of energy required to gasify liquid fuels versus solid fuels. The value of  $\alpha$ -cellulose has a relatively higher B-number than the other fuels because cellulose is the combustible component in corrugated cardboard and wood, while other substances in the fuels char and slow the combustion process. Additionally, the value for  $\alpha$ -cellulose shown in the figure is an ideal value that does not incorporate any losses, which results in a dramatic shift in its B-number value. In general, a lower B-number indicates a higher pyrolysis temperature because the fuel requires more energy to gasify. Therefore, a larger B-number indicates a fuel that has a higher thermodynamic efficiency during combustion [29].

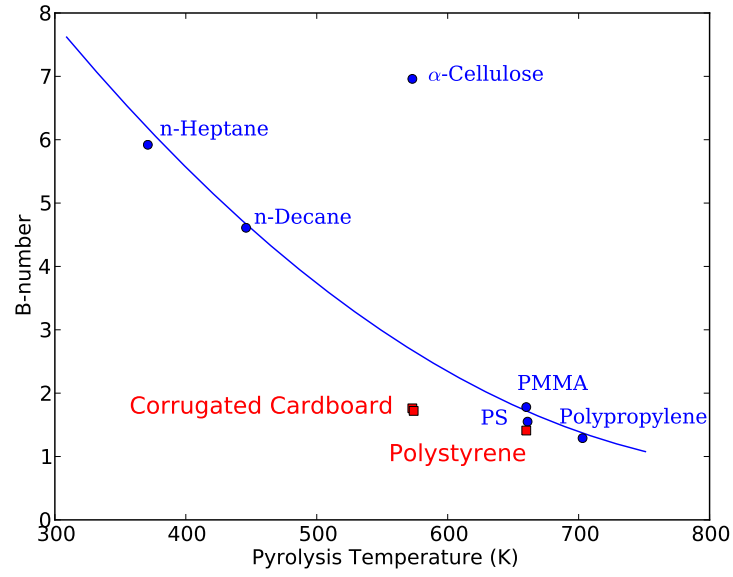


Figure 2: Values of the  $B$ -number for a range of fuels. The circles are  $B$ -number values from Annamalai and Sibulkin [27] that were calculated using only the thermodynamic properties. The red squares show the  $B$ -number values that were obtained experimentally by this study.

A simple description of mass transfer is derived by rearranging Eq. 1 for the mass-loss rate to yield the expression

$$\dot{m}_f'' = \frac{\bar{h}}{c_p} \ln(1 + B) \quad (3)$$

with the assumption of a unity Lewis number in which the thermal and mass diffusivities are assumed to be equal. Equation 3 will be used to determine a B-number for a given fuel by experimentally measuring the mass-loss rate.

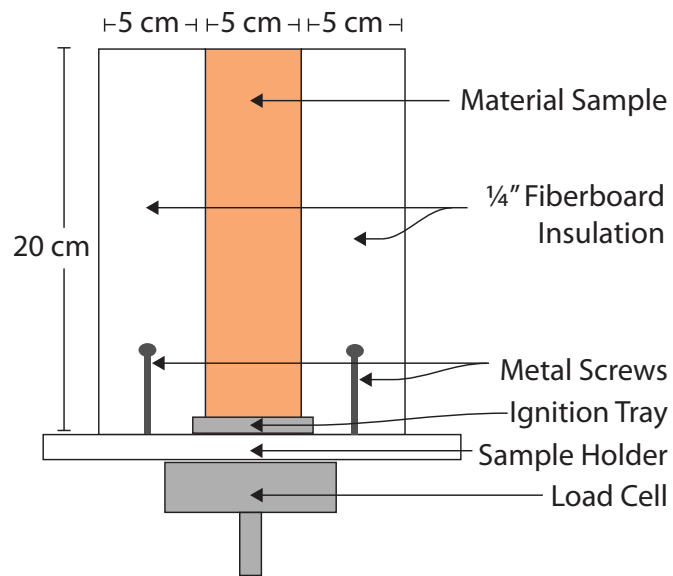
As explained in a recent publication [30], a number of improved laminar boundary layer types of theories result in formulas that are more complicated than Eq. 3, but the results are qualitatively the same. In larger tests that were previously performed, the fluctuating flames and the incipient turbulence raise questions about the degree of applicability of such theorems. For these reasons, this simple description of the mass transfer, Eq. 3, was chosen in this study over other relevant expressions.

In this study, the B-number is primarily a function of the material properties of a given fuel and it is obtained in a controlled experimental environment by assuming that the primary mode of heat transfer at the bench-scale is convection [31]. This assumption is reasonable for the small, laminar flames observed in this study. In examining Eq. 2, the B-number can be considered to be a ratio of the available energy (heat of combustion) to the energy required to gasify a given fuel (heat of gasification). Thus, the B-number is intrinsic to the properties of a material and is therefore independent of a particular scale. This allows for the results from the bench-scale tests to be used as a material input (instead of the heat release rate) for the prediction of large-scale warehouse fire behavior, and the effects of radiation at the large scale can be separately accounted for and will be further discussed in Section 5.

### 3. Experimental Setup and Procedure

Figure 3 shows a schematic of the experimental setup. A total of 13 tests were conducted using three different samples: single-wall corrugated cardboard (6 tests), polystyrene (3 tests), and single-wall corrugated cardboard backed with polystyrene (4 tests). The samples measured 5 cm wide by 20 cm in height; this aspect ratio was chosen because laminar flame spread was the primary focus of this study, and upwardly-spreading flames typically become turbulent above 20 cm [32]. For the bench-scale tests, a transition to a

## Front View



## Top View

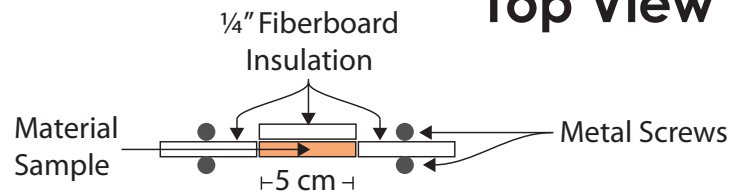


Figure 3: *Schematic of the experimental setup.*

turbulent regime was not considered for simplicity, which agreed with visual observations. For this study, the sample width was fixed at 5 cm to minimize the amount of variance between the tests and because a smaller sample size may affect the amount of combustible gases generated by the fuel due to significant diffusion of the fuel to the sides of the sample [33].

The typical mass of the samples was 4 g for corrugated cardboard and 36 g for polystyrene. Corrugated cardboard and polystyrene were chosen to be tested because they are the components of a Group A plastic commodity that is used to represent a worst-case fire scenario in large-scale warehouse tests. Additionally, corrugated cardboard is typically the first item to ignite and sustain flame spread in a warehouse fire. The mixed sample that consisted of corrugated cardboard backed by polystyrene was chosen as it is representative of a packed or mixed commodity that consists of both the corrugated cardboard packaging and the polystyrene material contained within. In this study, the results for the mixed sample were inconclusive thus far and are discussed in more detail later, but the inclusion of mixed fuel samples for the preliminary tests allows for the consideration of a mixed fuel in the presented framework. The measured quantities for each test included the mass-loss rate, flame height, and pyrolysis height.

The corrugated cardboard used in these tests was identical to the configuration and thickness that is used to package standard Group A plastics, and of the same type used in the small-scale tests that were performed by the authors in Part I [2]. The corrugated cardboard samples were of a type ‘C’ flute with a nominal thickness of 4 mm and 135 flutes per meter width [34] as shown in Figure 4(a). All of the tests were performed with the flutes aligned vertically along the 20 cm dimension, which is similar to the orientation of the flutes in an upright commodity box. The polystyrene samples were 3 mm thick as shown in Figure 4(b).

The mode of ignition for the tests was a small aluminum tray measuring 5 x 0.5 x 0.5 cm (Figure 3) that was placed at the base of the sample and contained a thin strip of glass fiber insulation soaked with n-heptane. This ensured a uniform mode of flaming ignition along the base of the fuel sample. The corrugated cardboard tests used 0.25 mL of n-heptane for ignition, whereas the polystyrene tests used 0.75 mL of n-heptane because it took a longer time for the polystyrene samples to ignite.

All of the fuel samples were insulated on the back and sides with 0.64 mm (0.25 inch) thick fiberboard insulation to isolate the burning to the front face of the samples only. The samples were secured in place by the insulating

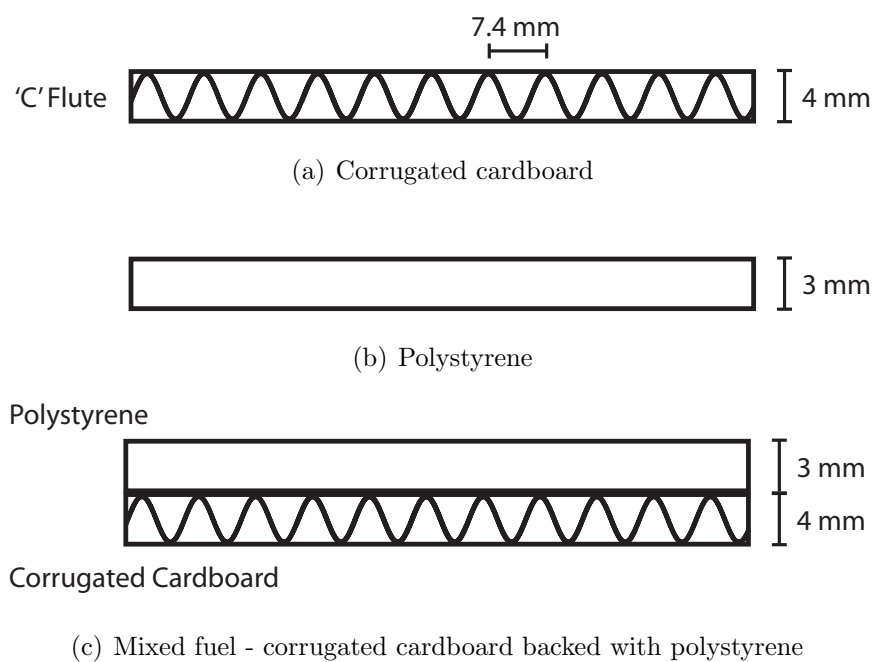


Figure 4: *Cross-sectional detail of the samples that were used in the tests: (a) Corrugated cardboard. The paper sheets are of a 26-26-26 lb. basis weight. (b) 4-mm thick polystyrene sheets. (c) Mixed-commodity sample (corrugated cardboard backed with polystyrene).*

fiberboard sheets that were supported by four metal screws attached to the 1.9 cm (0.75 inch) thick fiberboard base (Figure 3). All of the corrugated cardboard tests burned to completion and self-extinguished once the fuel was depleted. The polystyrene samples were manually extinguished after the flame reached a pyrolysis height of about 10 cm due to excessive smoke production and dripping on the bench-scale apparatus. However, the dripping and deformation of the polystyrene was not considered to be significant during the time frame considered in the results because the sample size in the experiment was small, and a significant accumulation of melted polystyrene was not observed during this time period.

The mass lost by the specimen was measured continuously using a load cell (Automatic Timing & Controls model 6005D) with an accuracy of  $\pm 0.5$  g as specified by the manufacturer. This is approximately 12% of the nominal initial mass of the corrugated cardboard samples and 2% of the nominal initial mass of the polystyrene samples. The load cell was calibrated prior to each test series using standard test weights. To measure the flame heights and record the burning history of the tests, video and still images were captured using a Sony Handycam HRR-SR5 model camera and a Canon EOS-5D digital single-lens reflex (DSLR) camera. Figure 5(a) depicts a visual time history of the vertical flame spread along a corrugated cardboard sample, and Figure 5(b) depicts the vertical flame spread along a polystyrene sample. The images were then loaded onto a computer, and a MATLAB image processing script was used to visually determine the flame heights as a function of time from each test. The flame height was defined as the tip of an attached yellow flame and was selected visually (by selecting the flame tip with the mouse pointer) from each picture by using the script. The processed images and resulting flame heights were consistent with visual comparisons from the test videos.

Similar to the flame heights, observations of the visual charring on the corrugated cardboard was used to determine the location of the pyrolysis front. For the polystyrene samples, visual bubbling and charring from the video were used to determine the location of the pyrolysis front. The corrugated cardboard and polystyrene tests were fairly repeatable, and the heights of the pyrolysis front in the laminar regime were fairly similar; thus, a best-fit functional approximation of the pyrolysis heights was made. This approximation was later used to determine an average mass-loss rate per unit area ( $\dot{m}_f''$ ), and finally, a B-number was calculated for each test. After the maximum pyrolysis height was reached, a constant height of 20 cm (for the corrugated



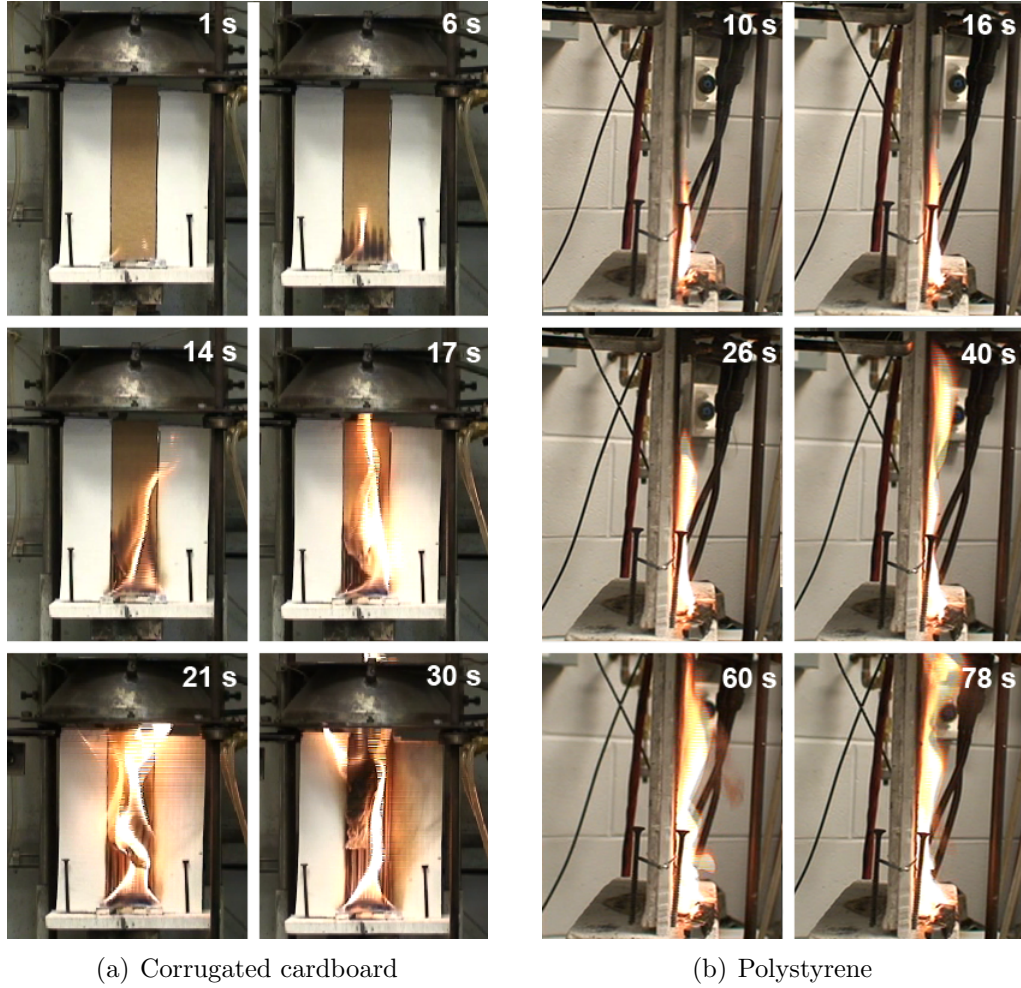


Figure 5: *Time history of the flame heights from the bench-scale tests. (a) Front view of the corrugated cardboard. (b) Side view of the polystyrene.*

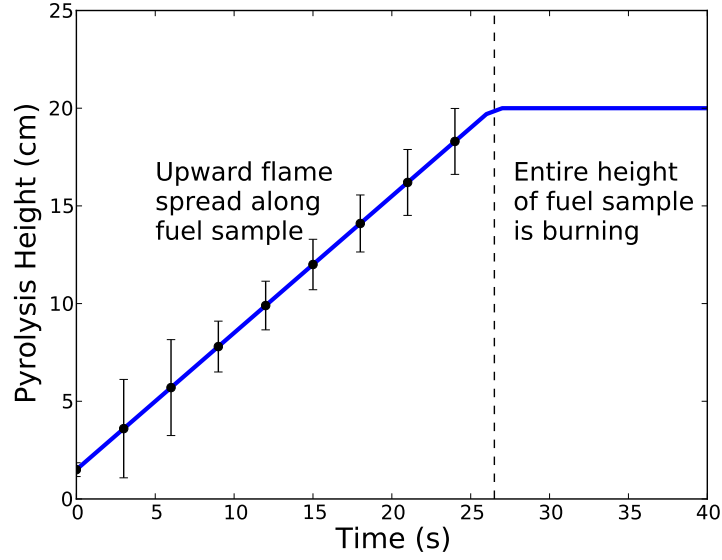
cardboard) or 10 cm (for the polystyrene) was assumed, which represents the entire surface of the front face of the sample.

#### 4. Experimental Results and Discussion

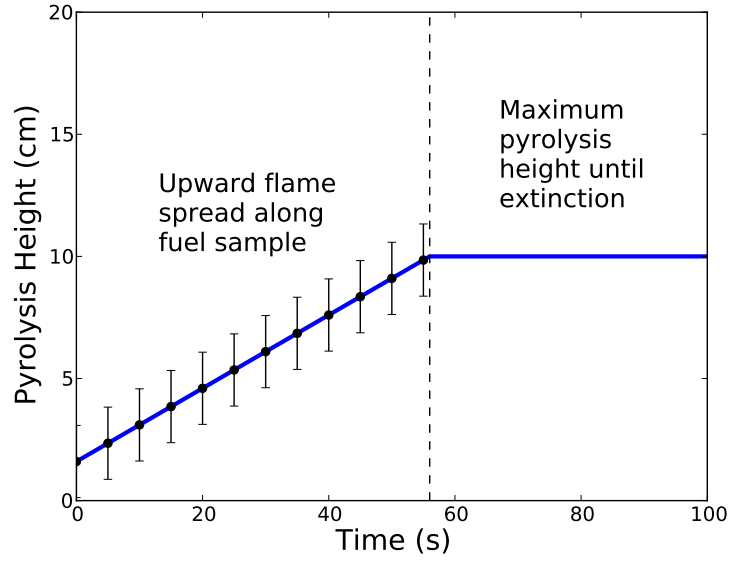
The results described in this section are based on a total of 13 bench-scale tests that were performed using the three samples that were discussed in Section 3. After uniformly igniting the base of the samples, the flame spread in the upward direction along the fuel samples. Due to edge effects along the fuel sample, a small amount of two-dimensional flame spread (both vertically and horizontally) occurred in the experiment. As the excess pyrolyzate burned above the pyrolysis zone, the unburned fuel above the pyrolysis zone ( $x_p$ ) was heated to its ignition temperature and the flame spread in the upward direction at an increasing rate [35]. As described in Section 3, the mass-loss rates were trimmed to only contain the time period during upward flame spread along the samples.

During the period of upward flame spread, the average value of  $\dot{m}_f''$  for corrugated cardboard was within a range of  $7.3 - 7.9 \cdot 10^{-4} \text{ g/cm}^2\text{-s}$ , and for polystyrene was within a range of  $6.7 - 6.8 \cdot 10^{-4} \text{ g/cm}^2\text{-s}$ . Figures 6(a) & 6(b) show the flame heights that were measured in the bench-scale experiments for corrugated cardboard and polystyrene and the pyrolysis height data fit that was used to determine the B-numbers for the corrugated cardboard and polystyrene samples. Using an average value from all of the tests that were performed on a given material sample, the B-number for corrugated cardboard was calculated to be 1.7 (standard deviation of 0.08) and for polystyrene was calculated as 1.4 (standard deviation of 0.02). The B-numbers were then input into the flame spread model as described in Section 5 to predict the flame heights for both the bench-scale and large-scale cases.

The mass-loss rate data were trimmed to contain only the time period where upward flame spread occurred along the sample. By reviewing the video recordings and mass-loss data for a particular test, Figure 7 shows the period after ignition and the omitted period after the flame reached the top of the sample. The trimmed (shaded) portion of the data was used to determine an average  $\dot{m}_f''$ . After the mass-loss rate was trimmed, it was then fit with a 4th-order polynomial to obtain a smooth mass-loss curve; the 4th order fits exhibited at least a 99%  $R^2$  value for each of the mass-loss data sets. The mass-loss data were then divided by the pyrolysis height data fits and the width of the sample to obtain an average  $\dot{m}_f''$  for each of the



(a) Corrugated cardboard



(b) Polystyrene

Figure 6: Flame heights as measured in the bench-scale tests. The points shows the measured flame heights with error bars representing the experimental range. The solid lines show the data fits for the flame height data, which was used to calculate the burning area.

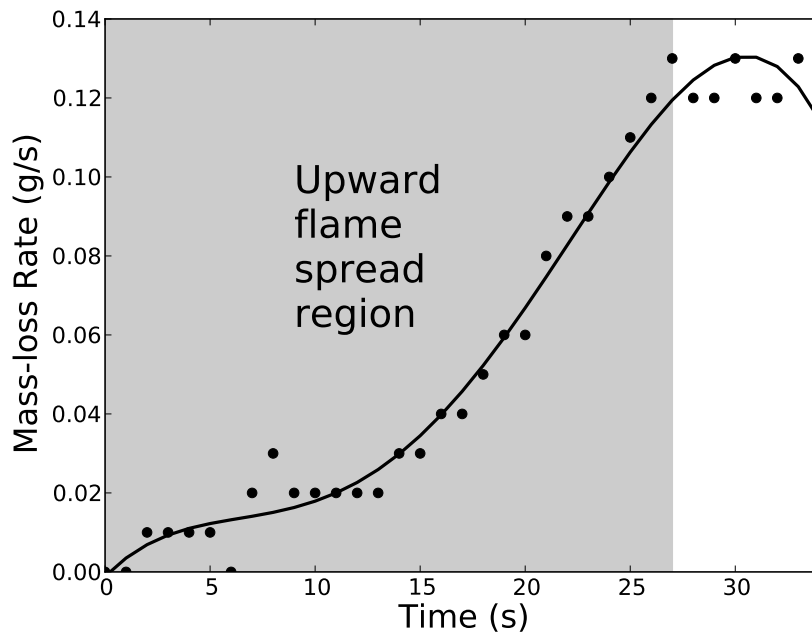


Figure 7: *Mass-loss rate for a representative corrugated cardboard test. The circles show the mass-loss data from the load cell, and the solid line shows the 4th-order data fit used to calculate the B-number. The shaded region shows the trimmed portion of the mass-loss rate during upward flame spread that was used to calculate the B-number.*

Table 2: Average  $\dot{m}_f''$  for each of the cardboard and polystyrene tests.

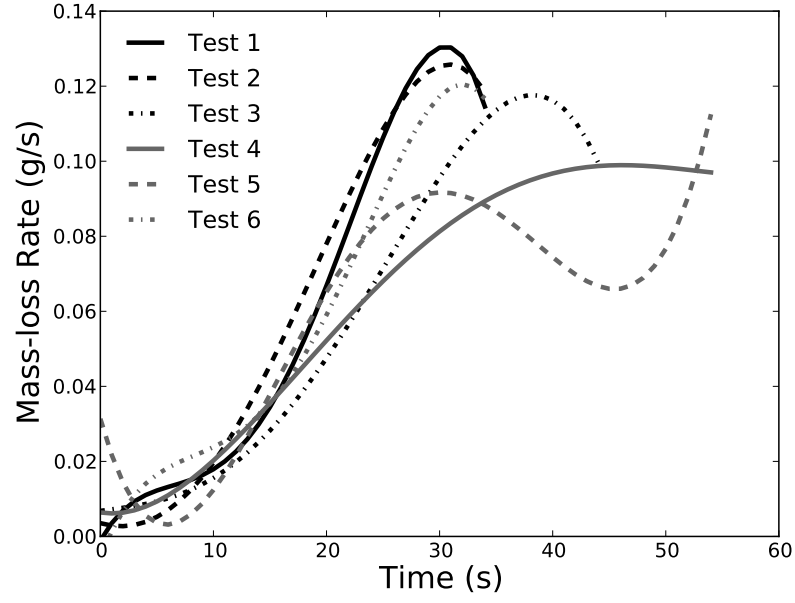
Test	$\dot{m}_f''$ (kg/m <sup>2</sup> -s)
Cardboard 1	7.7
Cardboard 2	7.8
Cardboard 3	7.4
Cardboard 4	7.5
Cardboard 5	7.9
Cardboard 6	7.3
Polystyrene 1	6.8
Polystyrene 2	6.7
Polystyrene 3	6.8

tests. The average  $\dot{m}_f''$  values for each of the cardboard and polystyrene tests are shown in Table 2. Figure 8 shows the mass-loss rates for each of the corrugated cardboard and polystyrene tests. In this figure, the initial time ( $t = 0$ ) corresponds to the time at which the sample was ignited and could sustain a flame without the pilot flame present. During this initial ignition period (typically less than a few seconds), some gasification occurred but was not sufficient to sustain the combustion of the material, which is the reason that some of the mass-loss rate fits exhibit a mass-loss rate slightly above zero at  $t = 0$ .

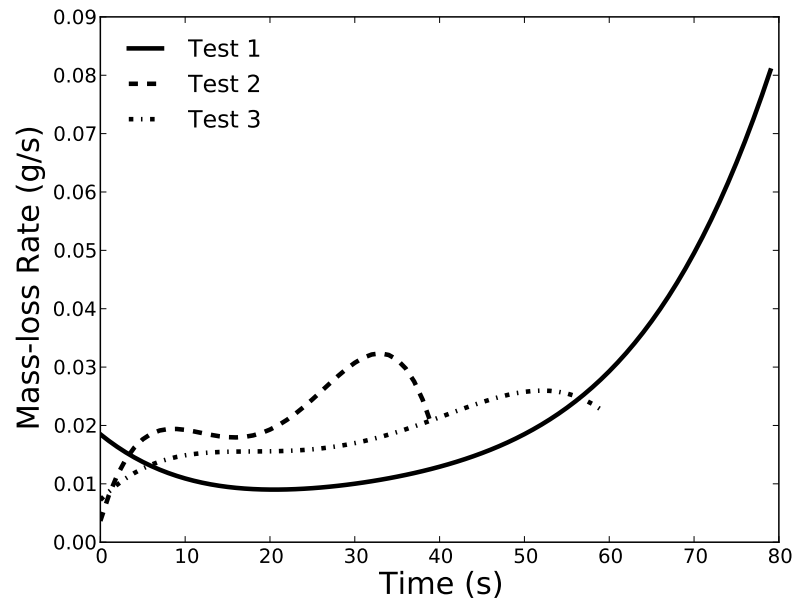
The value of  $\dot{m}_f''$  was then input into the experimental B-number formulation given by

$$B = \left( \frac{\dot{m}_f''}{\rho_g \alpha_g 0.59 / x_p [g x_p \beta \Delta T / \alpha_g \nu_g]^{1/4}} \right) - 1, \quad (4)$$

which uses a correlation for natural, laminar convection along a vertical plate,  $Nu = 0.59(GrPr)^{1/4}$  [34],  $\dot{m}_f''$  is the average mass-loss rate,  $\rho_g$  is the density of air,  $\alpha_g$  is the thermal diffusivity of air,  $x_p$  is the pyrolysis height,  $g$  is the acceleration due to gravity,  $\beta$  is the thermal expansion coefficient given by  $1/T_m$ ,  $\Delta T = T_m - T_\infty$ , and  $\nu_g$  is the kinematic viscosity of air. A mean gas temperature ( $T_m$ ) was used in the calculations by averaging the temperature of ambient gas,  $T_\infty = 20^\circ\text{C}$ , and an approximate flame temperature for cellulosic materials,  $T_f = 800^\circ\text{C}$  [36]. An equation similar to Eq. 4 for turbulent flow is derived fully in Part I of this paper [2]. Table 3 lists all of the values



(a) Corrugated cardboard



(b) Polystyrene

Figure 8: Mass-loss rate vs. time for the corrugated cardboard and polystyrene test samples.

Table 3: *Properties used in calculating the B-number (Eq. 4) estimated at a mean temperature of 683K [38].*

Property	Value
$\rho_g$	$0.50 \text{ kg/m}^3$
$\alpha_g$	$98 \cdot 10^{-6} \text{ m}^2/\text{s}$
Pr	0.7
$T_m$	683 K
$T_\infty$	298 K

used in Eq. 4. The thermo-physical properties of air are estimated at a mean gas temperature ( $T_m$ ) and are assumed to be constant [37].

## 5. Flame Spread Model

Figure 9 shows a schematic of the upward flame spread model from Sibulkin and Kim [39], which is a robust flame spread model that utilizes the B-number as an input parameter along with other thermophysical properties of the fuel, that was used to predict the flame heights at both the bench scale and the large scale. The pyrolysis zone is defined as the region of the solid fuel up to the pyrolysis height ( $x_p$ ) where combustible fuel vapors are outgassing. Some of the fuel burns directly in front of the combusting fuel surface, while some of the fuel is carried by buoyancy above its height of origin and burns above, which heats the virgin material in the preheat zone ( $\delta = x_f - x_p$ ) up to its ignition temperature. The fuel carried above the pyrolysis zone has been called *excess pyrolyzate* [35] and forms the physical flame height ( $x_f$ ) in which the resulting heat output drives the flame spread process. The rate of upward flame spread depends both on the amount of energy released by the combusting fuel and the rate at which the material pyrolyzes due to the flame heat flux,  $\dot{q}''(x)$ . This energy feedback from the gas phase to the condensed phase is the driving mechanism for the flame spread process. The B-number describes this feedback process as a nondimensional ratio.

The analytical model from Sibulkin and Kim [39] was adapted and solved numerically by using heat flux profiles from previous correlations. The heat flux is assumed to be constant along the pyrolysis region up to the pyrolysis height, and the flame spread occurs in one-dimension (vertically) along the

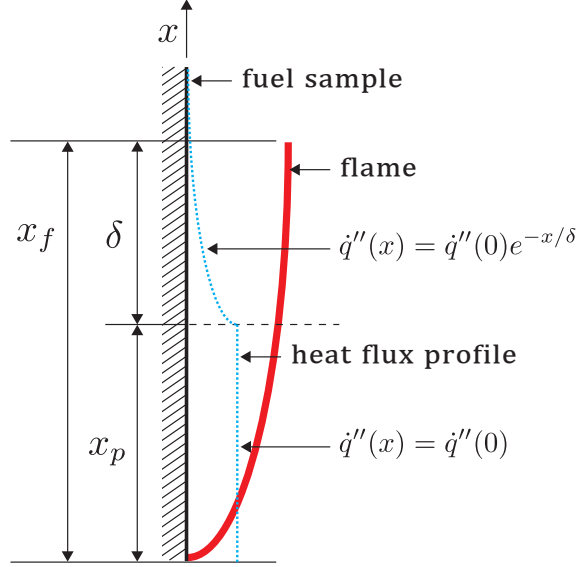


Figure 9: *The upward flame spread model proposed by Sibulkin and Kim [39].  $x_p$  is the pyrolysis height,  $x_f$  is the flame height,  $\delta$  is the preheat distance, and  $\dot{q}''(x)$  is the flame heat flux (dotted line).*

sample. In the preheat region ( $\delta$ ), the heat flux decays exponentially as a function of distance ( $x$ ), which follows from the heat flux distribution measurements by Sibulkin and Lee [40]. This heat flux condition is detailed in Eqs. 5a & 5b. Once the material in the preheat region reaches its pyrolysis temperature, it begins to outgas combustible vapors and the pyrolysis region grows, resulting in a larger flame height and more energy feedback to the unburned fuel; then the process repeats. Therefore, the process of upward flame spread can be thought of as a moving ignition front, similar to the leapfrogging process first described by de Ris [41].

Following from the concept of a moving ignition front, when solving the flame spread process numerically, the material sample is discretized into sections measuring 0.1 cm in height, and the initial conditions for the pyrolysis height and flame height that were observed in the experiments at the ignition time ( $t=0$ ) are input into the model. The heat flux profile is mapped along the height of the sample (as in Figure 9) by applying the following boundary conditions at the surface

$$\dot{q}''(x) = \dot{q}''(0) \exp(-x/\delta) \quad \text{for } x > x_p \quad (5a)$$



Table 4: *Physical properties used in the flame spread model for corrugated cardboard and polystyrene*

Property	Corrugated Cardboard	Polystyrene	Units
$k$	0.06 [42]	0.12 [27]	$W/m-K$
$\rho_s$	8.39 [43]	1.07 [27]	$g/m^3$
$c_p$	1.20 [43]	1.34 [27]	$J/g-K$
$\Delta H_c$	14,090 [44]	23,610 [44]	$J/g$
$\Delta H_g$	2,200 [43]	1,590 [43]	$J/g$
$T_p$	573 [27]	660 [27]	$K$

$$\dot{q}''(x) = \dot{q}''(0) \quad \text{for } x \leq x_p \quad (5b)$$

where  $\dot{q}''(0)$  is constant but can be modified to account for radiation from the flame,  $x$  is the height along the fuel sample, and  $\delta$  is the preheat region. A Nusselt number correlation is used to describe the turbulent, convective heat transfer process that is present at the large scale. An appropriate correlation was used for natural, turbulent convection along a vertical plate,  $Nu = 0.13(GrPr)^{1/3}$  [34]. Using this definition for the convective heat transfer coefficient, the initial heat flux,  $\dot{q}''(0)$ , to be used in Eqs. 5a & 5b can be approximated as

$$\dot{q}''(0) = \dot{q}_c'' = \bar{h}_c(T_f - T_\infty), \quad (6)$$

which neglects radiant heat transfer, and where  $\bar{h}_c$  is the convective heat transfer coefficient,  $T_f$  is the flame temperature for cellulosic materials ( $T_f = 800^\circ\text{C}$ ) [36, 37], and  $T_\infty$  is the ambient temperature. This results in a total heat flux from the flame,  $\dot{q}''(0)$ , of  $5.2 \text{ kW/m}^2$  for this configuration. Heat fluxes that incorporate both convection and radiation will be later discussed when considering large-scale warehouse radiation effects, which essentially modifies the  $\dot{q}''(0)$  term in Eqs. 5a & 5b.

After the heat flux is mapped along the height of the sample for the first time step, the forward heating parameter,  $\phi$ , is calculated and later used to find the velocity of the pyrolysis front. The forward heating parameter ( $\phi$ ) was defined by Sibulkin and Kim [39] as the ratio of the forward heat transfer rate to the rate of heat release per unit width of the fuel ( $\phi = \dot{q}_F' / \dot{q}_c'$ ). The forward heat transfer rate ( $\dot{q}_F'$ ) is calculated by the integral of the heat flux

above the pyrolysis length ( $x_p$ ) as in

$$\dot{q}'_F = \int_{x_p}^{\infty} \dot{q}''(x) dx, \quad (7)$$

where  $\dot{q}''(x)$  is the heat flux along the height of the sample (Eq. 5a), and  $\dot{q}'_c$  is the rate of heat release per unit width of the sample given by  $\dot{q}'_c = \dot{m}'_f \Delta H_c$ . An expression for the mass flux from the pyrolysis region ( $\dot{m}'_f$ ) obtained by Sibulkin and Kim and used in their flame spread model [39] is given by

$$\dot{m}'_f(x_p) = 0.59 \frac{\mu_f}{Pr^{3/4}} \left( \frac{g\beta\Delta T}{\nu_g} \right)^{1/4} \ln(1+B)x_p^{3/4} \text{ (laminar)}, \quad (8a)$$

$$\dot{m}'_f(x_p) = 0.13 \frac{\mu_f}{Pr^{2/3}} \left( \frac{g\beta\Delta T}{\nu_g} \right)^{1/3} \ln(1+B)x_p^{3/4} \text{ (turbulent)}, \quad (8b)$$

where  $\mu_f$  is the viscosity of air,  $Pr$  is the Prandtl number,  $g$  is the acceleration due to gravity,  $\beta$  is the thermal expansion coefficient,  $\Delta T$  is defined as  $(T_m - T_\infty)$ ,  $\nu_g$  is the kinematic viscosity of air,  $B$  is the B-number for the material as calculated by Eq. 4, and  $x_p$  is the pyrolysis height. The flame spread model switches to the turbulent formulation if the flame height ( $x_f$ ) becomes greater than 20 cm in length [32], which is later used when validating the model against large-scale fire test data. Once the forward heating parameter ( $\phi$ ) is calculated from  $\phi = \dot{q}'_F/\dot{q}'_c$ , the velocity of the moving pyrolysis front for the current time step is calculated by

$$V(x_p) = A_l \phi x_p^{1/2} \text{ (laminar)}, \quad (9a)$$

$$V(x_p) = A_t \phi x_p \text{ (turbulent)}, \quad (9b)$$

where the terms  $A_l$  and  $A_t$  are given in by Sibulkin and Kim [39] by

$$A_l = \frac{\Delta H_c \Delta H_g}{(4/3)\rho_s c_s k_s (T_p - T_\infty)^2} \left[ 0.59 \frac{\mu_f}{Pr^{3/4}} \left( \frac{g\beta_f \Delta T}{\nu_g^2} \right)^{1/4} \ln(1+B) \right]^2 \text{ (laminar)}, \quad (10a)$$

$$A_t = \frac{\Delta H_c \Delta H_g}{\rho_s c_s k_s (T_p - T_\infty)^2} \left[ 0.13 \frac{\mu_f}{Pr^{2/3}} \left( \frac{g\beta_f \Delta T}{\nu_g^2} \right)^{1/3} \ln(1+B) \right]^2 \text{ (turbulent)}, \quad (10b)$$

where  $\Delta H_c$  is the heat of combustion,  $\Delta H_g$  is the heat of gasification,  $\rho_s$ ,  $c_s$ , and  $k_s$  are thermophysical properties of the condensed phase material,  $T_p$  is the pyrolysis temperature of the condensed phase material, and the remaining terms were defined in Eqs. 8a & 8b. Table 4 lists the condensed phase material properties for corrugated cardboard and polystyrene that are used in Eqs. 10a & 10b.

The resulting change in the flame height (due to the upward velocity of the pyrolysis front) is added to the current pyrolysis height for the next time step as  $x_p[t + \Delta t] = x_p[t] + V[t] \cdot dt$ . In the final calculation of the time step, the pyrolysis height is converted to the height of the flame tip by using an expression by Annamalai and Sibulkin [26] for natural convection as given by

$$x_f = 0.64(r/B)^{-2/3}x_p. \quad (11)$$

where  $r$  is the mass consumption number given by  $(Y_{O,\infty}/\nu_s)$ , and  $B$  is the B-number for the material. The values of  $r$  were used as 0.194 for cardboard and 0.0749 for polystyrene [27]. This assumption of a constant ratio of the flame height to the pyrolysis height is based on the simplification that the burning rate is a function of the incident heat flux and that all of the excess fuel above the pyrolysis region is burned [26, 35]. After the new flame height ( $x_f$ ) is calculated, the numerical routine continues to the next time step and the process repeats starting from Eqs. 5a & 5b. This results in the prediction of the flame height as a function of time, i.e., a flame spread prediction.

To predict the flame heights in large-scale warehouse fires, both convection and radiation are incorporated into the flame spread model, which effectively modifies the  $\dot{q}''(0)$  term in Eqs. 5a & 5b. The simplest method for incorporating radiation is to use the Stefan-Boltzmann equation [34] to represent the radiant heat transfer from the gas phase by adding a radiative component to Eq. 6, resulting in

$$\dot{q}''(0) = \dot{q}_c'' + \dot{q}_r'' = \bar{h}_c(T_f - T_\infty) + \epsilon\sigma(T_f^4 - T_\infty^4), \quad (12)$$

where  $\dot{q}_c''$  is the convective heat flux,  $\dot{q}_r''$  is the radiative heat flux,  $\epsilon$  is the emissivity of the fuel assumed to be unity, and  $\sigma$  is the Stefan-Boltzmann constant ( $5.67 \cdot 10^{-8} \text{ W/m}^2 \cdot \text{K}^4$ ). This results in a total flame heat flux in which  $\dot{q}''(0) = 13 \text{ kW/m}^2$ .

A more detailed and useful method for representing the radiation in a large-scale warehouse setting is to incorporate a radiant heat-flux correlation based on work by de Ris and Orloff [11] for radiant heat transfer between

parallel panels. As previously mentioned, in this study, the flow and geometry effects were separated from the effects of the mass transfer. Therefore, this expression is useful for various geometries, including the geometry in a warehouse storage configuration, in which the flame spread can be considered to be occurring between two parallel plates represented by the flue space between the rows of stored commodities. Following the method of de Ris and Orloff [11], the radiant heat flux is calculated using

$$\dot{q}_r'' = \left( \frac{\varsigma_p \dot{q}_A''' w^2 d}{2x_f w} \right) + \dot{q}_{loss}'', \quad (13a)$$

$$\text{and } \varsigma_p = \frac{\beta_1(Y_s + Y_g)^{1/4} \varsigma_f}{\varsigma_f + \alpha_p \varsigma_f + \alpha_p} - \frac{2\beta_2 \dot{q}_{loss}''}{d \dot{q}_A'''} \quad (13b)$$

where  $\varsigma_p$  is the nondimensional panel width,  $\dot{q}_A'''$  is the volumetric heat release rate assumed to be  $1110 \text{ kW/m}^3$  [11],  $w$  is the sample width,  $d$  is the separation distance of the panels,  $x_f$  is the flame height,  $\dot{q}_{loss}''$  is the surface heat loss rate fixed at a constant value of  $5 \text{ kW/m}^2$  [11], and  $\beta_1$  and  $\beta_2$  are constants equal to 1.04 and 1.7, respectively.  $Y_s$  is the soot yield of the fuel equal to 0.01 g/g for a low-sooting fuel such as corrugated cardboard, and  $Y_g$  is added to the soot yield to account for radiation from the combustion gases for fuels having little to no soot and is equal to 0.01 g/g [11].  $\varsigma_f$  is the nondimensional flame height equal to  $x_f/w$ , and  $\alpha_p$  is the aspect ratio equal to  $d/w$ .

In this formulation for the radiant heat flux, an increase in the panel separation distance ( $d$ ) results in an increased radiant heat flux because the space between the panels is assumed to be fully occupied by flames. Thus, the separation distance for this study was fixed at 0.15 m (6 inches), which is representative of the flue space that is present in a typical warehouse commodity fire test. This expression for the radiant heat flux is dependent on both the flame height and the soot yield of the fuel, which are important factors to consider when modeling flame spread at the warehouse scale. In this study, a representative value for the soot yield ( $Y_s$ ) was chosen as 0.01 g for a cellulosic material such as corrugated cardboard, which is a very low sooting fuel. The soot yields are assumed to be constant; however, by using more information on the smoke point of the fuel from the bench-scale experiments, a variable soot yield can also be implemented. Using the results from Eqs. 13a & 13b for the radiant heat flux, a final expression for the flame heat

flux is given by

$$\dot{q}''(0) = \bar{h}_c(T_f - T_\infty) + \left( \frac{\varsigma_p \dot{q}_A''' w^2 d}{2x_f w} \right) + \dot{q}_{loss}'' \quad (14)$$

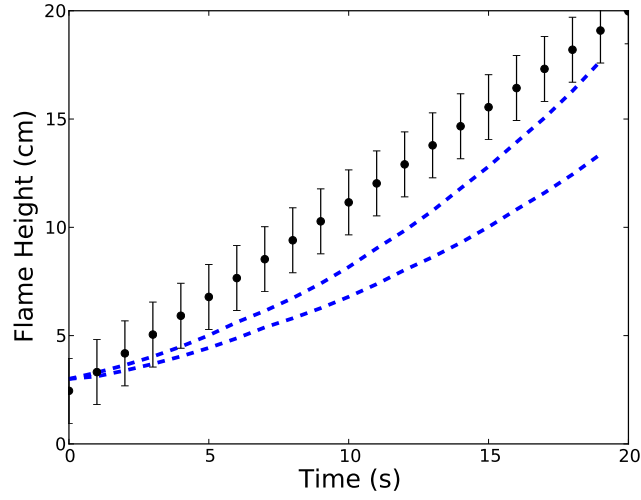
and results in a total flame heat flux in which  $\dot{q}''(0) = 27 \text{ kW/m}^2$ .

Quantifying the heat flux from the the flame to the fuel bed in complex geometries is a challenging yet important task that is required to further develop a quantitative description of warehouse fire behavior. Evaluating the radiative and convective heat flux fractions is not easily accomplished for a large assortment of practical geometries, but computational fluid dynamics software (CFD) allows for the possibility of modeling these complex flow conditions. If the pyrolysis rate of the fuels is effectively handled by the B-number in such CFD codes, then the other flow conditions may be more easily resolved, which highlights the potential applications of this work in the future.

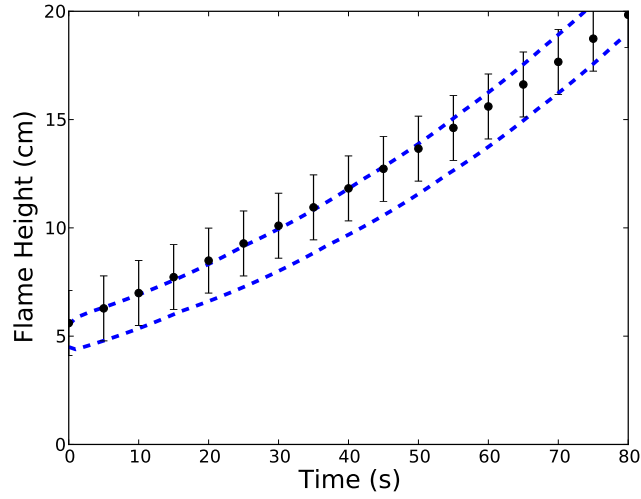
## 6. Flame Spread Model Results and Discussion

The results from the flame spread model were compared to the bench-scale results by using the observed flame heights from the videos of each of the 13 tests. Figure 10(a) shows the flame heights for corrugated cardboard as predicted by the model versus the bench-scale flame heights from the experiments. The flame height predictions for corrugated cardboard are in good agreement with the experimental flame heights. Figure 10(b) shows the flame heights for polystyrene as predicted by the model versus the bench-scale flame heights from the experiments. The flame height predictions for polystyrene are in good agreement with the experimental flame heights at the bench-scale. The bench-scale predictions are in reasonable agreement with the experimental flame heights because the dominant mode of heat transfer in the experiments was assumed to be laminar, natural convection on a vertical plate, and the same mode of heat transfer was used in the flame spread model. Under these conditions, the initial heat flux was calculated to be  $\dot{q}''(0) = 5.5 \text{ kW/m}^2$ . The thermal behavior of the fuel samples was considered to be a slab of finite thickness, and a more detailed analysis can be found in Overholt [45].

The results from the flame spread model were then compared to the large-scale by using the flame heights from the rack-storage warehouse fire tests.



(a) Corrugated cardboard



(b) Polystyrene

Figure 10: *Flame heights in the bench-scale tests compared to the predicted flame heights. The points show the measured flame heights with error bars that represent the amount of deviation between the tests. The dashed lines show the upper and lower bounds of the predicted flame heights by incorporating the standard deviation of the calculated  $B$ -number.*

The flame heights for the large-scale warehouse fires were obtained from video data from three large-scale warehouse commodity fire tests that were performed at Underwriter’s Laboratory in Northbrook, Illinois [46]. The fuel consisted of paper cups (Class III commodity) that were packed in corrugated cardboard boxes and stacked between 6.1 m to 9.1 m in height (20 and 30 ft) in a rack-storage configuration. The boxes were ignited along the bottom edge in the flue space between the racks. Flame heights as a function of time were acquired from the test videos. Figure 11 shows a snapshot from a warehouse fire test as the flame spreads up through the flue space between the boxes.

The average B-number for corrugated cardboard (1.7) was used in the large-scale flame spread predictions because it is nondimensional and describes the mass flux for both the bench-scale and large-scale scenarios. Previous studies have shown that the B-number is not constant, but varies to some degree in both time and space [24]. For the purposes of the large-scale flame height predictions, the B-number was assumed to have a constant value of 1.7.

The flame spread model predictions for flame height were validated for a range of experimental flame heights from the large-scale UL tests, and the results are shown in Figure 12. The shaded region shows the range of experimental flame heights from three large-scale tests as extracted from the test videos. The three dashed lines show the flame height predictions using the experimentally determined B-number for three different flame heat fluxes. To account for the various modes of heat transfer that are present in the large-scale, three different methods for representing the flame heat flux,  $\dot{q}''(0)$ , were used in the flame spread model as described in Section 5. Curve (a) used a flame heat flux equal to  $27 \text{ kW/m}^2$ , Curve (b) used a heat flux equal to  $13 \text{ kW/m}^2$ , and Curve (c) used a heat flux equal to  $5.2 \text{ kW/m}^2$ .

The flame heat flux that resulted in the best flame height predictions accounts for both convective and radiative heat transfer by using a radiation correlation based on heat transfer between two parallel plates as shown in Eq. 13a. This is the most representative of the fire conditions in the large-scale warehouse fire tests because the fire is ignited in the flue space between the commodity boxes and spreads upwards between the stack of commodity boxes. In this case, radiant energy feedback was occurring between the parallel fuel surfaces as the flames grew larger and increased the flame heat flux and the flame spread rate. The model shows good agreement for the initial stage of fire growth at the large-scale in which the primary fuel is the



Figure 11: *A large-scale fire test as the flame spreads up through the flue space between the packed commodity boxes [46].*

cardboard packaging of the cartons.

## 7. Conclusions

This work has developed a bench-scale method to experimentally determine the B-number to rank the flammability hazard of a given material. The results from the bench-scale tests were then used to model vertical flame spread in the flue space during a warehouse test with commodity stacked up to a height of 9.1 m (30 ft). The flame spread model that showed the best agreement with the large-scale experimental flame heights (Figure 12) used the flame heat flux that incorporates both convective heat transfer and a correlation for radiative heat transfer between parallel plates. Therefore, the processes of heat transfer (flow conditions) and mass transfer (B-number) were successfully decoupled and were expressed independently of one another, which enabled the scaling of the results from the bench-scale tests to the large-scale warehouse conditions. The B-number was obtained from bench-scale experiments where the flow conditions can be controlled to better understand the effects of material properties. Three different flow conditions were used to model heat transfer in the large-scale and validated by using large-scale commodity fire test data.

Additionally, because the soot yield ( $Y_s$ ) is nondimensional and intrinsic to a given material, it can be a useful parameter to model the radiation effects at the large-scale. As  $Y_s$  increases, the radiant feedback from the gas



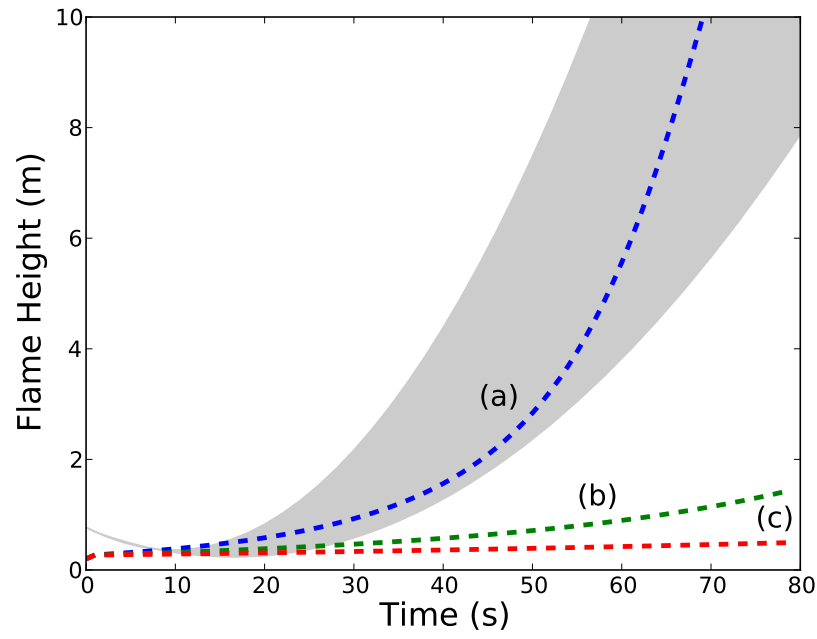


Figure 12: Flame heights from the large-scale UL experiments (gray band) are compared to the predicted flame heights (dashed lines) using three different heat flux models, which correspond to three values of the initial heat flux,  $\dot{q}''(0)$ : a)  $27 \text{ kW/m}^2$ , b)  $13 \text{ kW/m}^2$ , c)  $5.2 \text{ kW/m}^2$ . The gray band shows the range of flame heights as measured from the large-scale experiments, and the dashed line shows the predicted flame heights.

phase combustion to the fuel increases, which results in an increased rate of flame spread. The B-number was also calculated for a sample of corrugated cardboard backed by polystyrene, which is representative of a mixed commodity. As a first estimation of the influence of commodity within corrugated cardboard packaging, the mass-loss rates for the mixed sample was discussed. This relates to the objective of the experimental method to determine a quantified flammability ranking for materials consisting of both homogeneous and mixed commodities. Future work involves more understanding of the physical interaction between multiple material samples to quantify the effects of a mixed commodity on the overall flame spread process.

The two parameters, the B-number and the soot yield ( $Y_s$ ), can both be determined from bench-scale test methods and utilized in a flammability ranking scheme that is valid in large-scale fires. This is important because flammability ranking is coupled with the upward flame spread process, which is the most significant hazard in a warehouse storage fire, and the B-number and soot yield seem to describe the process well for the vertical flue space in the warehouse scenario. This study establishes the framework for a more cost-effective method to determine the flammability hazard of various commodity materials using a simple bench-scale test method.

A framework was demonstrated for which the results from bench-scale tests can be used to quantitatively rank the flammability of both single fuels and mixed commodity configurations and predict flame heights at the large-scale. If the pyrolysis rate of the fuels is effectively handled by the B-number in CFD codes, then the flow conditions for more complex geometries may be more easily resolved, which highlights the potential applications of this work in the future. The B-number and soot yield are fundamentally robust parameters that may be used in the future as means to classify the flammability of a given warehouse commodity, to strengthen the level of confidence in ranking a commodity, and to increase the effectiveness of warehouse fire protection and suppression applications. Additionally, the results of this study are useful for the application of sprinkler activation and determining the amount of sprinkler suppression that is necessary as a fire grows larger.

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